

$B(\tau^- \rightarrow \eta \pi^- \nu_\tau)$ , the  $a_0(980)$ ,  
and  
New Weak Interactions

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# Outline

- Motivation
- Estimate of  $B(\tau^- \rightarrow \eta \pi^- \nu_\tau)$ :
  - Vector contribution
  - Scalar contribution
- Sensitivity to new interactions
- Implications of various  $B(\tau^- \rightarrow \eta \pi^- \nu_\tau)$  values

# Motivation

- $\tau^- \rightarrow \eta \pi^- \nu_\tau$  is sensitive to new-physics
- CLEO limit on  $B(\tau^- \rightarrow \eta \pi^- \nu_\tau) < 1.4 \times 10^{-4}$  (90% CL)
- Chiral perturbation theory:  $B(\tau^- \rightarrow \eta \pi^- \nu_\tau) \approx 1.3 \times 10^{-5}$ 
  - Pich, PLB 196, 561 (1987); Tisserant and Truong, PLB 115, 264 (1982); Neufeld, and Rupertsberger, Z. Phys. C 68, 91 (1995).
  - Assumes  $a_0^-(980) \rightarrow \eta \pi^-$  contribution dominates, and that  $a_0(980)$  is a  $q\bar{q}$  state
  - However,  $a_0(980)$  is probably a 4-quark,  $u\bar{d}s\bar{s}$  state
- At this level,  $\rho^-(770) \rightarrow \eta \pi^-$  could also lead to a vector contribution
- Therefore, we...
  - calculate the  $\rho$  contribution to  $\tau^- \rightarrow \eta \pi^- \nu_\tau$
  - perform an alternative calculation of the  $a_0$  contribution
  - comment on sensitivity of  $B(\tau^- \rightarrow \eta \pi^- \nu_\tau)$  to new weak interactions

# Vector contribution


- Note:  $B(\tau^- \rightarrow \pi^- \pi^0 \nu_\tau)$  is large (25.5%) and completely dominated by  $\rho^-$  contribution
- So expect  $\rho^-$  to also dominate the vector contribution to  $\tau^- \rightarrow \eta \pi^- \nu_\tau$ , with branching fraction

1 power from phase space,  
2 from vector amplitude

$$B_{L=1}(\tau^- \rightarrow \eta \pi^- \nu) = \left( \frac{g_{\rho\eta\pi}}{g_{\rho\pi\pi}} \right)^2 \left( \frac{P_{\rho \rightarrow \eta\pi}}{P_{\rho \rightarrow \pi\pi}} \right)^3 B(\tau^- \rightarrow \rho^- \nu)$$

# Obtaining $g_{\rho\eta\pi}$ Coupling

- $\rho^- \rightarrow \eta\pi^-$  has not been observed, but we can obtain  $g_{\rho\eta\pi}$  from the Dalitz-plot distribution of  $\eta \rightarrow \pi^+\pi^-\pi^0$ 
  - Method used by Ametller & Bramon, PRD 24, 1325 (1981)
  - Now more precise data, access to more terms in Dalitz-plot distribution
- Assume the decay has 3 contributions: scalar,  $\rho^+$  and  $\rho^-$ :
  - ( $\rho^0$  is forbidden due to C conservation)

$$M_{\eta \rightarrow \pi^+\pi^-\pi^0} \equiv M_{+-0} = M_S + M_{\rho^+} + M_{\rho^-}$$


- Scalar part has flat distribution in the  $\pi^+\pi^-\pi^0$  Dalitz plot, and is also the only contribution to  $\eta \rightarrow 3\pi^0$

$$|\mathcal{M}_S|^2 = 8(2\pi)^3 m_\eta \Gamma_\eta \mathcal{B}(\eta \rightarrow \pi^0 \pi^0 \pi^0) \frac{6\sqrt{3}}{Q^2 S_1} \frac{3!}{9}$$

$$= 0.065,$$

$$Q \equiv m_\eta - 3m_\pi$$

$$S_1 \equiv \int dX dY = 2.75 = \text{area of Dalitz plot}$$

Dalitz plot variables

$$X = \frac{\sqrt{3}(T_{\pi^+} - T_{\pi^-})}{Q}, \quad Y = \frac{3(T_{\pi^0} - 1)}{Q}$$

Write vector part as:

$$M_{\rho^\mp} = -\underline{g_{\rho\eta\pi}} g_{\rho\pi\pi} \frac{(P_\eta + P_{\pi^\pm}) \cdot (P_{\pi^\mp} - P_{\pi^0})}{(P_{\pi^\mp} + P_{\pi^0})^2 - m_\rho^2 + i\Gamma_\rho m_\rho}$$

The coupling we are after

- Expand squared amplitude to 3<sup>rd</sup> order in

$$r = \frac{m_\eta Q}{m_\rho^2 - m_\eta^2 / 3 - m_\pi^2 - i\Gamma_\rho m_\rho} = 0.14 - 0.03i$$

- Obtain total rate relative to scalar part:

$$|M_{+-0}|^2 \propto 1 + \alpha Y + \beta Y^2 + \gamma X^2 + \delta(Y^3 + YX^2)$$

- Where

$$\alpha = -4 \underline{g_{\eta\rho\pi} g_{\rho\pi\pi}} \Re\{\mathcal{M}_S^* r\} \frac{1}{|\mathcal{M}_S|^2}$$

$$\beta = \left[ -\frac{4}{3} g_{\eta\rho\pi} g_{\rho\pi\pi} \Re\{\mathcal{M}_S^* r^2\} + 4(g_{\eta\rho\pi} g_{\rho\pi\pi})^2 |r|^2 \right] \frac{1}{|\mathcal{M}_S|^2}$$

$$\gamma = \frac{4}{3} g_{\eta\rho\pi} g_{\rho\pi\pi} \Re\{\mathcal{M}_S^* r^2\} \frac{1}{|\mathcal{M}_S|^2}$$

$$\delta = \left[ \frac{4}{9} g_{\eta\rho\pi} g_{\rho\pi\pi} \Re\{\mathcal{M}_S^* r^3\} + \frac{8}{3} (g_{\eta\rho\pi} g_{\rho\pi\pi})^2 \Re\{r(r^2)^*\} \right] \frac{1}{|\mathcal{M}_S|^2}$$

- Dalitz-plot parameters measured by KLOE (arXiv:0707.2355):

$$|M_{+-0}|^2 \propto 1 - 1.09Y + 0.124Y^2 + 0.057X^2 + 0.14Y^3 + ?YX^2$$

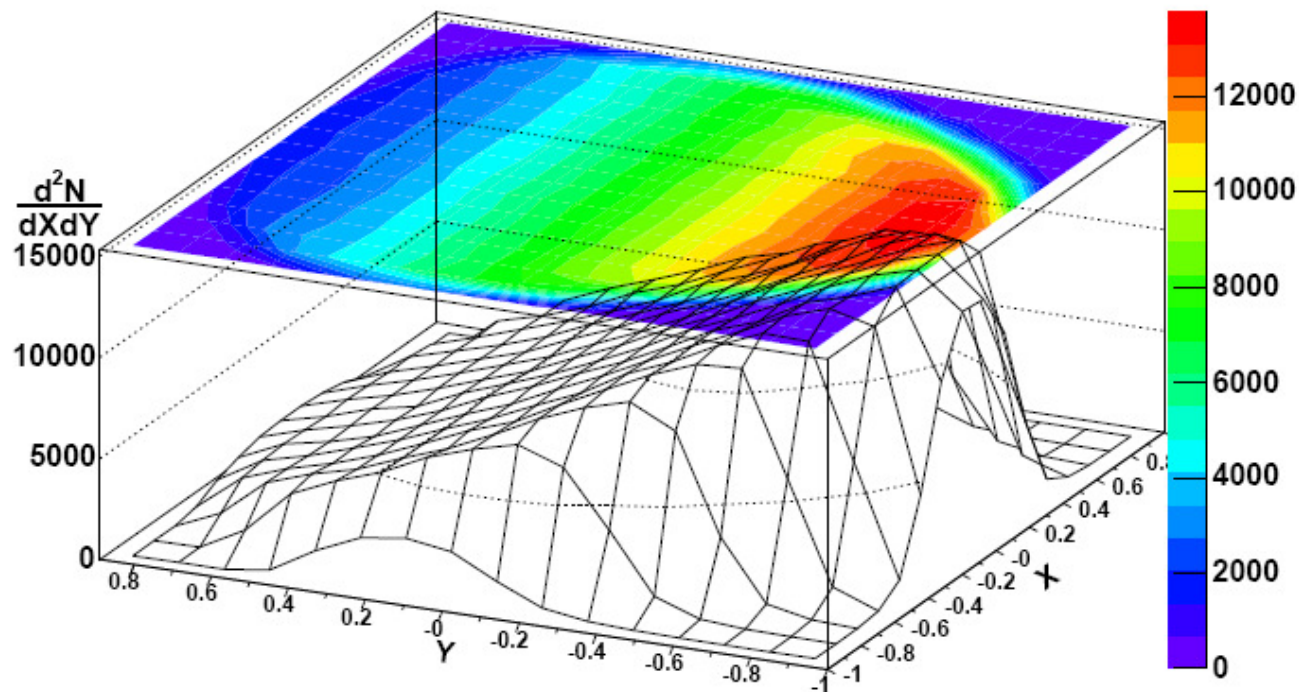


Fig. 2. Dalitz-plot distribution for the whole data sample. The plot contains 1.34 millions of events in 256 bins.



- Comparing the Y coefficients, get:

$$\underline{g_{\eta\rho\pi}g_{\rho\pi\pi}} = \frac{1.09}{4} \frac{\mathcal{M}_S}{\Re(r)} = 0.51$$

Taken to be real

- Extract  $g_{\rho\pi\pi}$  from  $\rho \rightarrow \pi\pi$  width:

$$g_{\rho\pi\pi} = \sqrt{\frac{6\pi m_\rho^2 \Gamma_\rho}{P_{\rho \rightarrow \pi\pi}^3}} = 6.0.$$

$$\underline{g_{\eta\rho\pi}} \approx 0.085$$

Consistent w. Ametller & Bramon

- So the vector contribution to  $B(\tau \rightarrow \eta\pi\nu)$  is

$$B_{L=1} = \left( \frac{g_{\rho\eta\pi}}{g_{\rho\pi\pi}} \right)^2 \left( \frac{P_{\rho \rightarrow \eta\pi}}{P_{\rho \rightarrow \pi\pi}} \right)^3 B(\tau^- \rightarrow \rho^- \nu) \approx 3.6 \times 10^{-6}$$

# Cross-checks

- The other coefficients are a test of the model:

$$|M_{+-0}|^2 \propto 1 - 1.09Y + 0.27Y^2 + 0.05X^2 + 0.03(Y^3 - YX^2)$$

- Compare with KLOE measurement:

$$|M_{+-0}|^2 \propto 1 - 1.09Y + 0.124Y^2 + 0.057X^2 + 0.14Y^3$$

- Floating  $\arg(M_S) = 15^\circ$  improves agreement only slightly
- Also check ratio of BR's::

$$\frac{B(\eta \rightarrow \pi^+ \pi^- \pi^0)}{B(\eta \rightarrow 3\pi^0)} = \begin{cases} 0.7 & , \text{measured} \\ 0.71 & , \text{model + KLEO parameters} \\ 0.76 & , \text{model + our parameters} \end{cases}$$

# Scalar Contribution to $B(\tau^- \rightarrow \eta \pi^- \nu_\tau)$

- Chiral perturbation theory calculates assumed  $a_0(980)$  is a  $q\bar{q}$  state and are complicated
- We conduct a simpler estimate and arrive at a similar result:
- Vector current is conserved up to  $m_d - m_u$ :

$$\nabla^\mu \bar{u}(x) \gamma_\mu d(x) = (m_d - m_u) \bar{u}(x) d(x) + \text{EM term}$$

- We estimate the scalar matrix element by relating the P-wave states  $a_0(980)$  &  $a_1(1260)$ :

$$\frac{B_{L=0}(\tau^- \rightarrow a_0^-(980) \nu)}{B(\tau^- \rightarrow a_1^-(1260) \nu)} \sim 1.3 \frac{\left| \langle 0 | S | a_0(980) \rangle \right|^2}{\left| \langle 0 | A | a_1(1260) \rangle \right|^2} \left( \frac{m_d - m_u}{m_{a_1(1260)}} \right)^2$$

Phase space

$\sim 1$ , since fixed by quark-model wave functions

$$B_{L=0} \sim 10^{-5}$$

# Limits on New Physics

- $B(\tau^- \rightarrow \eta \pi^- \nu)$  can be used to put bounds on new scalar interactions up to the SM expectation  $B(\tau^- \rightarrow \eta \pi^- \nu) \sim 10^{-5}$
- A limit  $B(\tau^- \rightarrow \eta \pi^- \nu) < 3 \times 10^{-5}$  implies

$$\frac{M_{\text{Scalar}}}{M_W} > \left(3 \times 10^{-5}\right)^{-\frac{1}{4}} \sim 13$$

for the same couplings as in the SM.

- Competitive with limits from angular distributions in nuclear  $\beta$  decay,  $\sim 4$  (expected to improve to  $\sim 7$  and then to  $\sim 15$ )
- The two limits are complementary:
  - $\beta$ -decay: 1<sup>st</sup>-generation couplings
  - $\tau^- \rightarrow \eta \pi^- \nu$ : 3<sup>rd</sup>-generation couplings

# Conclusions

- Our estimates

- $B_{L=0}(\tau^- \rightarrow \eta \pi^- \nu) \sim 10^{-5}$
- $B_{L=1}(\tau^- \rightarrow \eta \pi^- \nu) \approx 3 \times 10^{-6}$

imply the following for the measured value of  $B_{L=0}(\tau^- \rightarrow \eta \pi^- \nu)$ :

- $\approx 3 \times 10^{-6}$ , especially with a  $\rho^-(770)$  peak  $\Rightarrow$ 
  - No surprises
- $\approx 10 \times 10^{-6}$ , especially with a  $a_0^-(980)$  peak  $\Rightarrow$ 
  - $a_0^-(980)$  is a  $q\bar{q}$  state after all
- $> \sim 30 \times 10^{-6}$ , especially with scalar dominance  $\Rightarrow$ 
  - Possibly new scalar interactions,  $M_S \sim 13 M_W$  for weak coupling
- Note that BaBar has limit  $B(\tau^- \rightarrow \eta' \pi^- \nu) < 7.2 \times 10^{-6}$ 
  - Contributions from additional intermediate resonances

# Backup Slides

# Arguments for 4-quark $a_0(980)$

- $\Gamma(a_0(980)) \sim 50 \text{ MeV}$ , compare with  $\Gamma(\rho(770)) \sim 150 \text{ MeV}$
- $B(a_0(980) \rightarrow K\bar{K}) / B(a_0(980) \rightarrow \eta\pi) = 18\%$   
despite highly limited phase space in  $KK$  final state
- $a_0(980)$  production much suppressed wrt.  $q\bar{q}$  states.
  - This suppression might be smaller in heavy ion collisions due to high quark multiplicity, especially  $s$  quarks

- 2<sup>nd</sup>-class currents: Weinberg, Phys. Rev. 112, 1375 (1958):
  - 1<sup>st</sup>-class currents:  $GJ_1G^{-1} = -\xi J_1$ ,  $JPG = 0^{++}, 0^{--}, 1^{+-}, 1^{-+}$
  - 2<sup>nd</sup>-class currents:  $GJ_2G^{-1} = +\xi J_2$ ,  $JPG = 0^{+-}, 0^{-+}, 1^{++}, 1^{--}$
  - where  $\xi = 1(-1)$  for S, A, P (V, T) and  $G = Ce^{i\pi I_2}$ :

- Full expression for  $\eta \rightarrow 3\pi^0$ :

$$|\mathcal{M}_S|^2 = 8(2\pi)^3 m_\eta \Gamma_\eta \mathcal{B}(\eta \rightarrow \pi^0 \pi^0 \pi^0) \frac{6\sqrt{3}}{Q^2 S_1} \frac{3!}{9}$$

$$= 0.065,$$

- Rho decay amplitude:  $\Gamma(\rho \rightarrow \pi\pi) = \frac{P_{\rho \rightarrow \pi\pi}}{32\pi^2 m_\rho^2} \int |M_\rho|^2 d\Omega$

$$M_\rho = g_{\eta\rho\pi} \varepsilon_\mu (P_{\pi^+} - P_{\pi^-})^\mu$$