

Role of axialvector mesons near the chiral phase transition

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in collaboration with

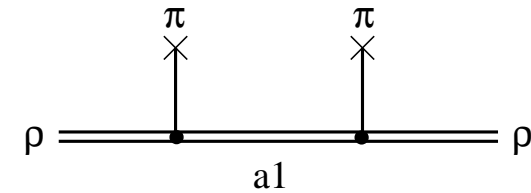
M. Harada (Nagoya) and W. Weise (TUM)

references:

- Harada and Sasaki, Phys. Rev. D **73**, 036001 (2006).
- Harada, Sasaki and Weise, arXiv:0805.4792 [hep-ph];
arXiv:0807.1417 [hep-ph], accepted for publication in Phys. Rev. D.

Chiral symmetry restoration: role of a_1

- V-A mixing in matter through thermal pions:



- at low temperature $T \ll m_\pi$:

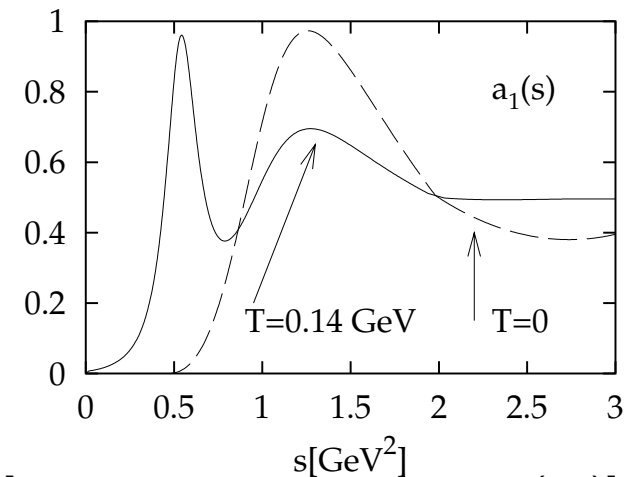
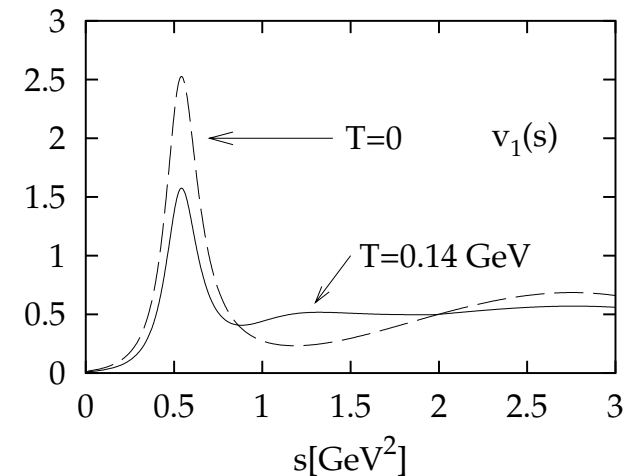
only pions are thermally activated.
a low-energy “theorem” as

$$G_V^{\mu\nu}(T) = (1 - \epsilon)G_V^{\mu\nu}(0) + \epsilon G_A^{\mu\nu}(0)$$

$$G_A^{\mu\nu}(T) = (1 - \epsilon)G_A^{\mu\nu}(0) + \epsilon G_V^{\mu\nu}(0)$$

$$\epsilon = \frac{T^2}{6F_\pi^2}, \quad G_{V,A}: \text{current correlators}$$

[Dey, Eletsky and Ioffe (90)]



[Marco, Hofmann and Weise (02)]

- chiral symmetry restoration
 - “standard” picture: ρ and a_1 chiral partners
 - \Rightarrow degenerate masses $m_{a_1} - m_\rho \rightarrow 0$
 - $G_V - G_A \rightarrow 0$ at $T \rightarrow T_c$ in any case
 - maximal mixing $\epsilon = \frac{1}{2}$
 - $\Rightarrow G_A = G_V = \frac{1}{2} \left(G_V^{(\text{vac})} + G_A^{(\text{vac})} \right)$
 - \Rightarrow chiral sym. restoration? no a_1 modification?
- go beyond
 - low-energy limit, other hadrons than pions, higher temperature
- this work
 - the first systematic study within an effective field theory

An effective field theory for a π - ρ - a_1 system

- generalized hidden local symmetry (GHLS)

[Bando, Kugo, Yamawaki (85,88); Bando, Fujiwara, Yamawaki (88); Kaiser, Meissner (90)]

- an extension of non-linear chiral Lagrangian

- ρ , a_1 introduced as gauge bosons of the redundant gauge symmetry

- chiral perturbation theory with GHLS

[Harada and Sasaki (06)]

- systematics: order-by-order renormalization

- Weinberg 1st & 2nd sum rules

$$F_\pi^2 + F_{a_1}^2 = F_\rho^2, \quad F_{a_1}^2 M_{a_1}^2 = F_\rho^2 M_\rho^2$$

are stable against quantum corrections.

- description of chiral sym. restoration

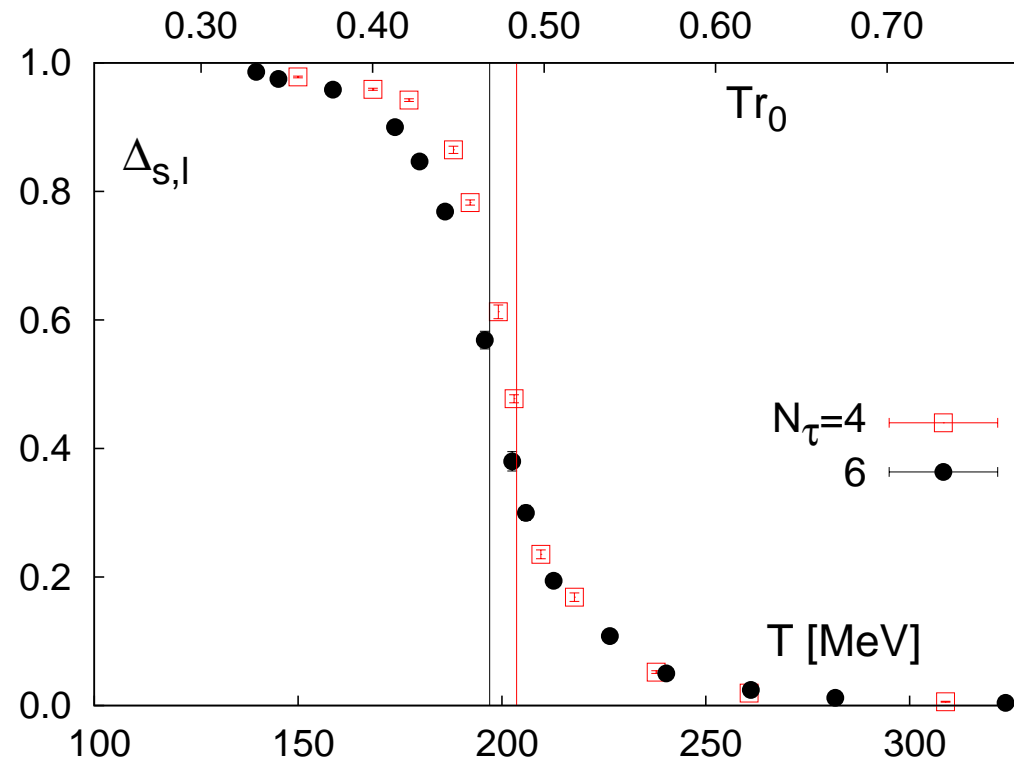
$$G_A - G_V \rightarrow 0 \quad \Rightarrow \quad M_{a_1} - M_\rho \rightarrow 0$$

just a parameter $\delta M = M_{a_1} - M_\rho$: how δM goes to zero?

Turning on/off partial restoration: flash temperature

- lattice QCD calculation with almost physical quark masses
($m_\pi \simeq 220$ MeV)

[M. Cheng *et al.* (08)]

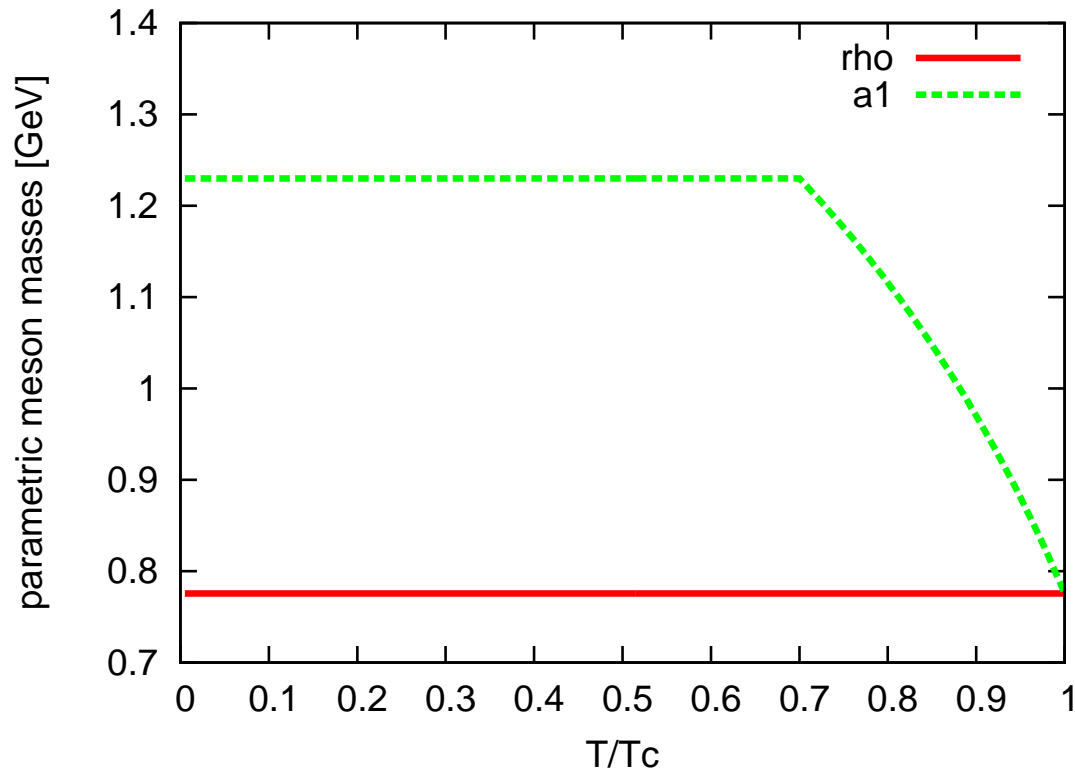


- quark condensate starts to melt at $T \simeq 140$ MeV
⇒ introduce the flash temperature $T_f = 140$ MeV:
partial chiral sym. restoration sets in at T_f .

- assume ρ bare mass is independent of T .
- take a simple T-dep. of $M_{a_1}^2 - M_\rho^2 = \delta M^2(T) = c(T)g^2 F^2$ as

$$c(T) = c(T = 0) \Theta(T_f - T) + c(T = 0) \Theta(T - T_f) \frac{T_c^2 - T^2}{T_c^2 - T_f^2}$$

- ρ and a_1 bare masses at finite T : ($T_f = 140$ MeV & $T_c = 200$ MeV)

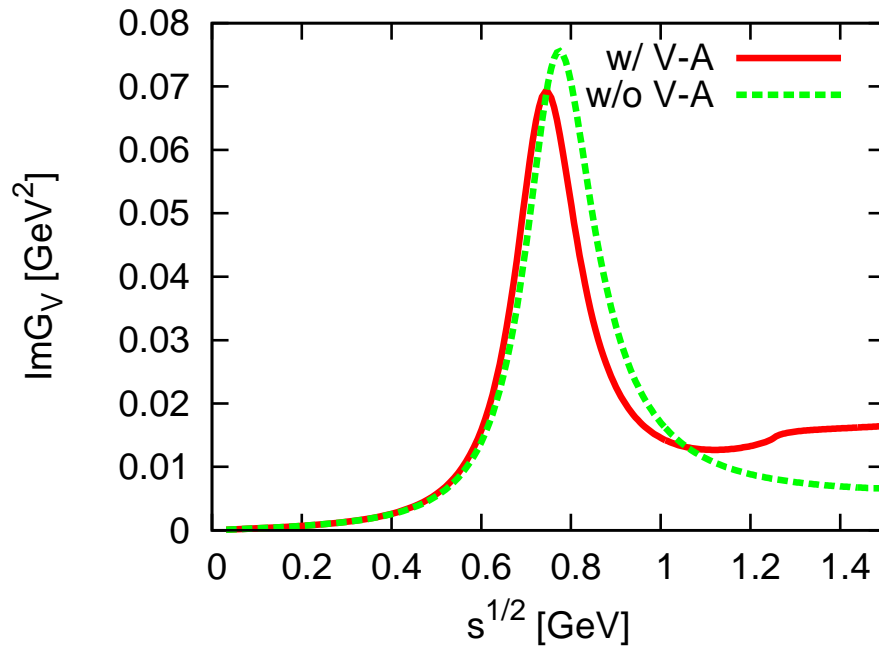


- pion decay constant $F_\pi^2 \propto c(T)F^2 \rightarrow 0$ at T_c

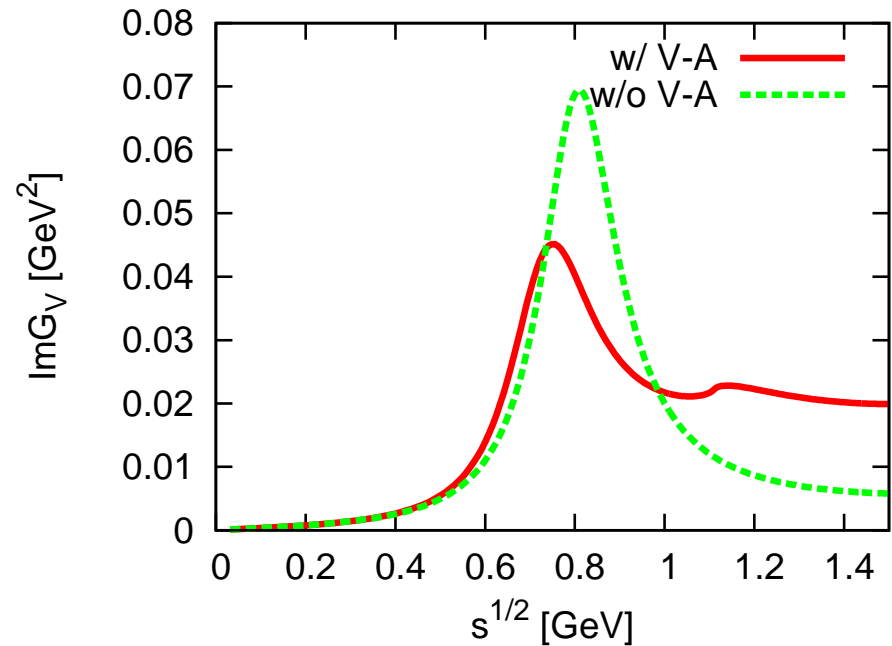
V-A mixing in hot matter: vector spectral function

- vector spectral function in the chiral limit

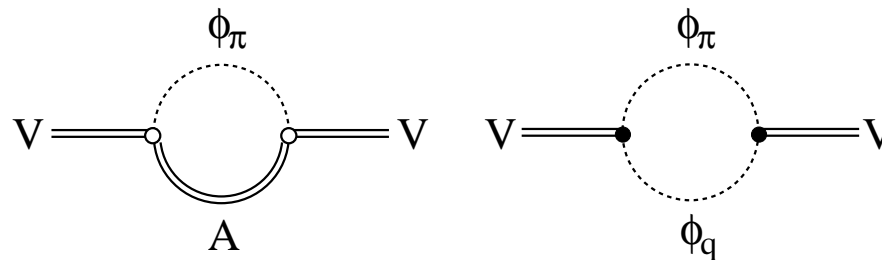
below T_f



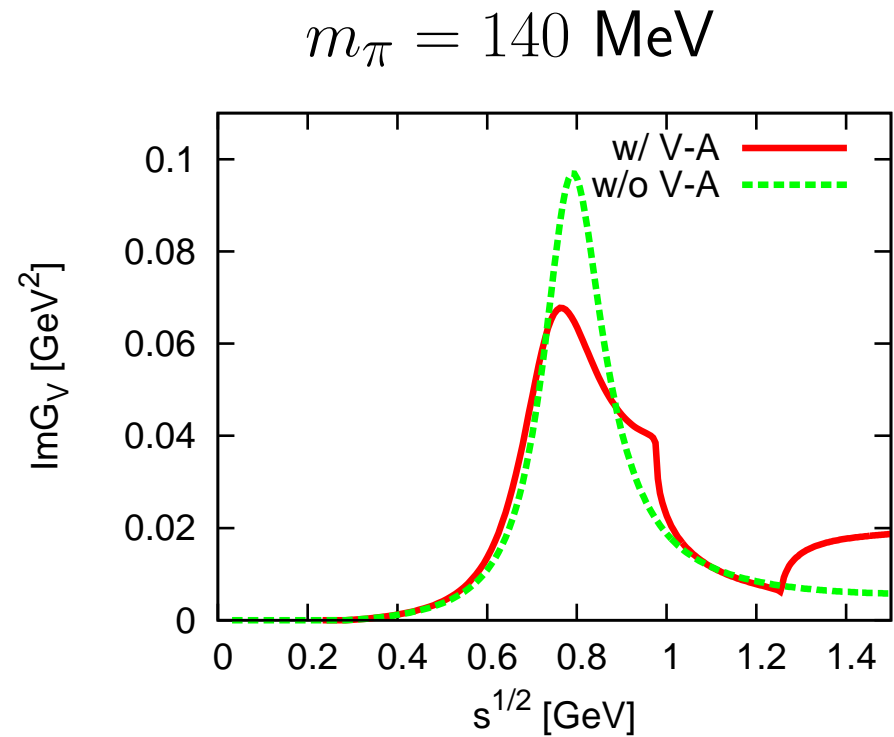
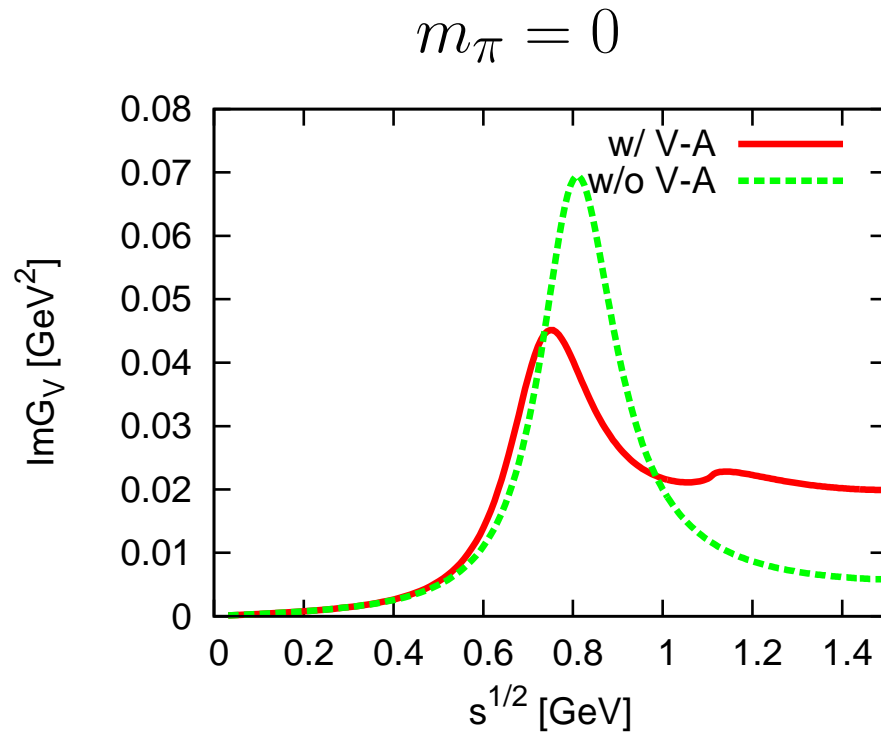
above T_f



- effect of a_1 meson becomes significant above T_f
- 2 bumps appear through V-A mixing

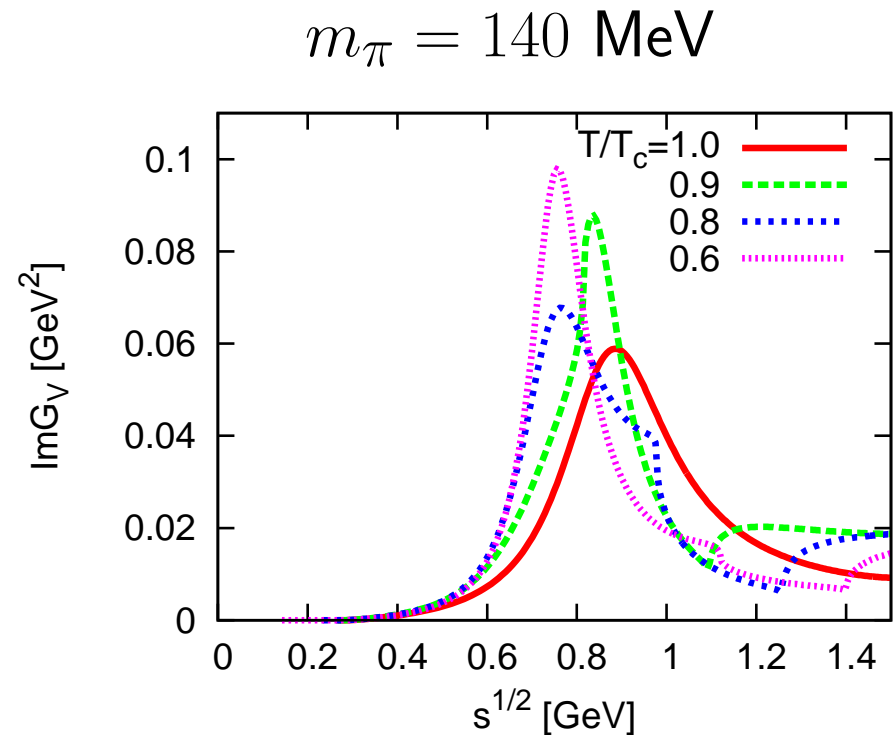
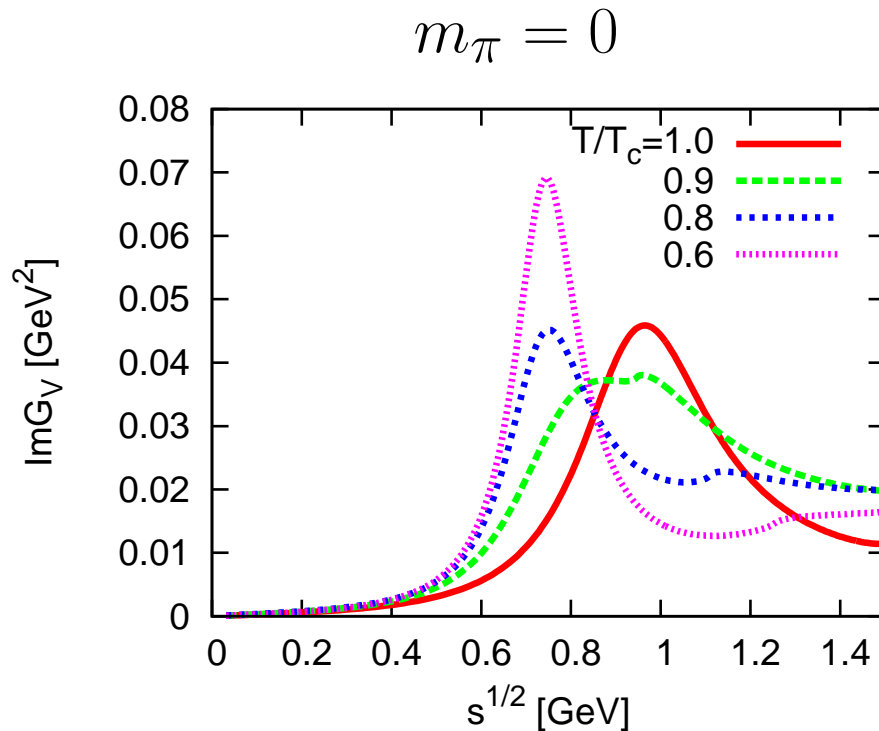


- vector spectral function: m_π dependence



- 2 processes kinematically allowed: $\rho + \pi \rightarrow a_1$ and $\rho \rightarrow a_1 + \pi$
- finite m_π : $\sqrt{s} = m_{a_1} - m_\pi$ and $\sqrt{s} = m_{a_1} + m_\pi$

- vector spectral function: T dependence



- 2 bumps on top at T_c
- masses of chiral partners near T_c :

$$M_{a_1} - M_\rho \simeq \sqrt{h_A - h_V} m_\pi \simeq 3 \text{ MeV}$$

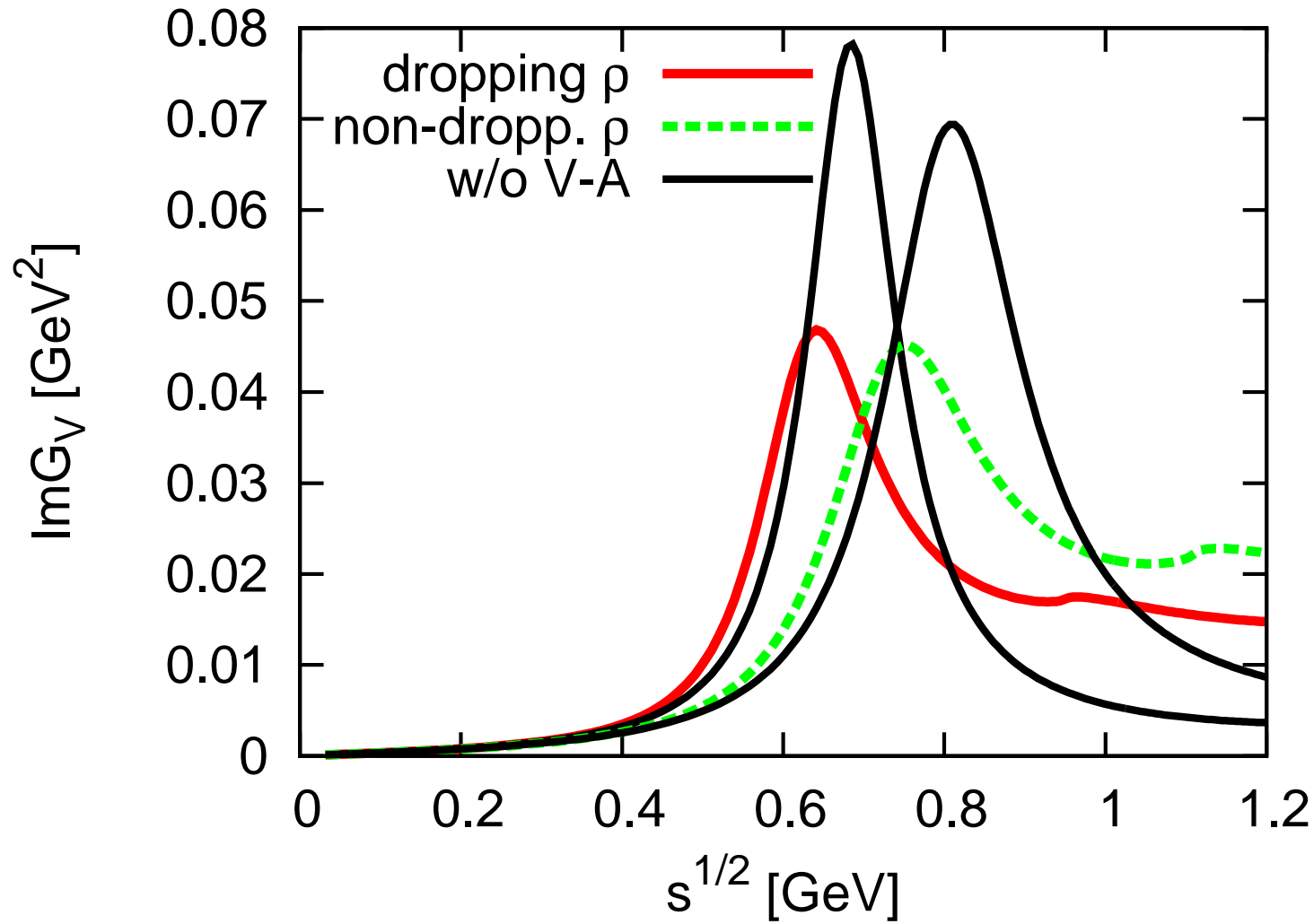
- unchanged by $\bar{q}q$ scalar mode

a linearized hidden-local-symmetric model

[M. Harada, C. S., W. Weise (2008)]

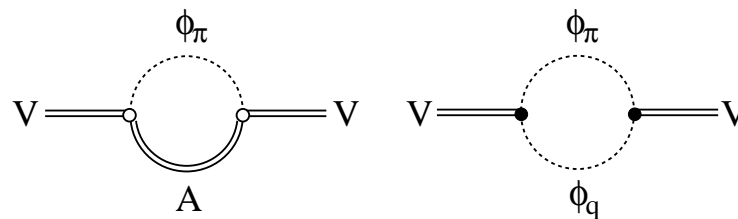
A case study: dropping vs. non-dropping ρ mass

- non-dropping vs. dropping *bare* ρ mass: V-A mixing more significant



Conclusions

- a systematic study of chiral symmetry restoration (d.o.f.: π, ρ, a_1)
 - vanishing order parameter: $F_\pi^2 \propto c(T)$
 - degenerate masses of chiral partners: $\delta M^2 \propto c(T)$
- single unified bump at T_c
 - \Leftrightarrow 2 bumps from a naive extrapolation of the mixing theorem
- ρ - a_1 masses well degenerate near T_c even for $m_\pi = 140$ MeV
 - suppression of a_1 - ρ - π coupling near T_c : $g_{a_1\rho\pi} \sim 0.06 m_\pi$



- pion loop contribution more and more important above T_f

$$\text{Im}\Pi_V(T) \sim \left(1 + \frac{M_\rho^2}{M_{a_1}^2(T)} \right) F(m_\pi, m_\pi; T)$$

- application to dilepton production, signature of χ -sym. restoration